SCIENTIFIC REPORTS

OPEN

Probing the structural evolution of ruthenium doped germanium clusters: Photoelectron spectroscopy and density functional theory calculations

Yuanyuan Jin[1,2], Shengjie Lu[3], Andreas Hermann[4], Xiaoyu Kuang[5], Chuanzhao Zhang[1], Cheng Lu[2,6], Hongguang Xu[3] & Weijun Zheng[3]

We present a combined experimental and theoretical study of ruthenium doped germanium clusters, $RuGe_n^-$ ($n = 3–12$), and their corresponding neutral species. Photoelectron spectra of $RuGe_n^-$ clusters are measured at 266 nm. The vertical detachment energies (VDEs) and adiabatic detachment energies (ADEs) are obtained. Unbiased CALYPSO structure searches confirm the low-lying structures of anionic and neutral ruthenium doped germanium clusters in the size range of $3 \leq n \leq 12$. Subsequent geometry optimizations using density functional theory (DFT) at PW91/LANL2DZ level are carried out to determine the relative stability and electronic properties of ruthenium doped germanium clusters. It is found that most of the anionic and neutral clusters have very similar global features. Although the global minimum structures of the anionic and neutral clusters are different, their respective geometries are observed as the low-lying isomers in either case. In addition, for $n > 8$, the Ru atom in $RuGe_n^{-/0}$ clusters is absorbed endohedrally in the Ge cage. The theoretically predicted vertical and adiabatic detachment energies are in good agreement with the experimental measurements. The excellent agreement between DFT calculations and experiment enables a comprehensive evaluation of the geometrical and electronic structures of ruthenium doped germanium clusters.

Cluster is a new state of aggregation, which can be regarded as the intermediate phases between atoms or molecules and bulk solids[1,2]. A systematic study of clusters can provide valuable information on the evolution of the structural and electronic properties when isolated atoms or molecules become larger agglomerations, and bridge many fields of physics, such as atomic, molecular, and condensed-matter physics[3]. The central issue in cluster science is the determination of the true global minimum structures[4]. However, as the size of cluster increasing, finding the global minimum geometries becomes increasingly difficult due to the much increased complexity of the potential surface as well as the rapid increase of the number of low-lying isomers[5,6]. Luckily, the combined experimental and theoretical photoelectron spectroscopy approach has become a very effective method to identify the cluster structures, which has successfully determined the true global-minimum structures for various clusters ranging from small or medium-sized cluster to large sized cluster[7–12].

Among the various classes of cluster materials, germanium-based clusters have attracted considerable experimental and theoretical attention because of that germanium is one of the most potential alternatives to silicon in microelectronic industry[13,14]. However, pure germanium clusters can not form stable fullerene-like cage

[1]Department of Physics and Optoelectronic Engineering, Yangtze University, Jingzhou 434023, China. [2]Department of Physics, Nanyang Normal University, Nanyang 473061, China. [3]State Key Laboratory of Molecular Reaction Dynamics, Institute of Chemistry, Chinese Academy of Sciences, Beijing 100190, China. [4]Centre for Science at Extreme Conditions and SUPA, School of Physics and Astronomy, The University of Edinburgh, Edinburgh EH9 3JZ, United Kingdom. [5]Institute of Atomic and Molecular Physics, Sichuan University, Chengdu 610065, China. [6]Department of Physics and High Pressure Science and Engineering Center, University of Nevada, Las Vegas, Nevada 89154, United States. Correspondence and requests for materials should be addressed to X.K. (email: scu_kuang@163.com) or C.L. (email: lucheng@calypso.cn) or H.X. (email: xuhong@iccas.ac.cn)

structures and are unsuitable as a building block of self-assembly materials[15,16]. In order to generate and stabilize Ge cage structures, numerous experimental and theoretical investigations have been performed on transition metal (TM)-doped germanium clusters, similar to the case of TM-doped silicon clusters[11,12,17–25]. It is found that the TM-doped germanium clusters reveal different growth pattern from the TM-doped silicon clusters[19,26]. The critical size of TM-encapsulated Si_n structures is generally suggested to appear at $n = 12$[10,27–29], while the TM-doped Ge_n clusters can form endohedral structures when $n \leq 10$[17–25]. Using solution chemistry methods, Wang et al.[17] successfully synthesized the intermetalloid $CoGe_{10}^{3-}$ cluster with a D_{5h} pentagonal prism structure. Zhou et al.[18] performed experimental measurements by using the standard Schlenk-line techniques and reported another pentagonal prismatic Zintl ion cage encapsulating an interstitial iron atom, $FeGe_{10}^{3-}$. Based on the anion photoelectron spectroscopy in combination with density functional theory (DFT) calculations, Deng et al.[11] studied the structural, electronic and magnetic properties of $VGe_n^{-/0}$ ($n = 3-12$) clusters and suggested that the endohedral structures occur from $n = 9$ and a D_{3d} distorted hexagonal prism cage structure is formed at $n = 12$. They also observed that the critical size of the transition from exo- to endohedral structures is $n = 9$ for both anionic and neutral $CoGe_n$ ($n = 2-11$) clusters[12]. Despite the much advances of $3d$ TM-doped germanium clusters, relatively little is known about the $4d$ TM-doped germanium clusters.

Very recently, Espinoza-Quintero et al.[30] successfully synthesized the 12-vertex endohedral cluster $RuGe_{12}^{3-}$, a previously unknown 3-connected polyhedral geometry of D_{2d}-symmetry, by the reaction of an ethylenediamine solution of K_4Ge_9 with $[Ru(COD)\{\eta^3\text{-}CH_3C\text{-}(CH_2)_2\}_2]$ (COD = 1,5-cyclooctadiene). In order to confirm the geometries and electronic properties observed in their experimental measurements, Goicoechea and McGrady[31] performed DFT calculations on MSi_{12} and MGe_{12} and concluded that the bicapped pentagonal prism structure of $RuGe_{12}^{3-}$ dominates the structural landscape for high valence electron counts (57–60). Nevertheless, up to now, there are no systematic investigations on neutral or single charged Ru-doped germanium clusters. There are still several essential open questions about the Ru-doped germanium clusters: How clusters grow with the increasing number of germanium atoms? What is the size of the smallest endohedral cage structures? What are the charge properties of the endohedral cage structures?

As an effort to address the above questions, here we report a combined photoelectron spectroscopy and DFT study on Ru-doped germanium clusters: $RuGe_n^-$ and $RuGe_n$ ($n = 3-12$). The vertical detachment energy (VDE) and adiabatic detachment energy (ADE) of $RuGe_n^-$ are estimated from their photoelectron spectra. The structures of $RuGe_n^-$ and $RuGe_n$ are assigned by the comparison of the theoretical simulations and experimental measurements. The thermodynamic stabilities of the obtained global minima are checked by analyzing the average binding energies (E_b) and the second energy difference ($\Delta^2 E$). The natural population analysis is conducted to trace the negative charge dense regions in the neutral and anionic forms.

Results and Discussion

Experimental results. The photoelectron spectra of $RuGe_n^-$ clusters from $n = 3$ to 12 measured at 266 nm photons are presented in Fig. 1. In each spectrum, the first peak represents the transition from the ground electronic state of the cluster anion to that of the corresponding neutral species, and the other peaks with higher binding energy denote transitions to excited electronic states of the neutral clusters. The VDE is evaluated from the first peak. Meanwhile, the ADE of each cluster is estimated by adding the instrumental resolution to the binding energy which is the interaction of the binding energy axis and a straight line drawn along the leading edge of the first peak. The experimental VDE and ADE values of $RuGe_n^-$ clusters from photoelectron spectra in Fig. 1 are summarized in Table 1.

In the spectrum of $RuGe_3^-$, the first peak occurs at around 2.08 eV, which gives the VDE. After this, there is a weak peak at 2.57 eV and then a broadened area with several obscure peaks between 3.00 and 4.00 eV. From the photoelectron spectrum of $RuGe_4^-$, the VDE is about 2.12 eV. The second peak at 2.52 eV is stronger than the first one whereas the third weak peak occurs at 2.90 eV. The spectrum of $RuGe_5^-$ reveals that the VDE is about 2.32 eV and a second smaller peak locates at 2.95 eV. For $RuGe_6^-$, the VDE is about 2.75 eV and the front three peaks of the photoelectron spectrum are very close, followed by a broadened area between 3.90 and 4.40 eV. In the case of $RuGe_7^-$, the spectrum exhibits a major peak at 2.53 eV, which is the VDE. For $RuGe_8^-$, the VDE is about 3.10 eV followed by a sharp peak at 3.62 eV. The photoelectron spectrum of $RuGe_9^-$ reveals the first weak peak at about 2.89 eV and a second small peak at 3.54 eV. The spectrum of $RuGe_{10}^-$ is similar to that of $RuGe_9^-$, in with a VDE is 3.17 eV and a second peak at 3.58 eV. In the case of $RuGe_{11}^-$, there is a broad area between 2.80 and 3.30 eV and the VDE is about 3.14 eV. For $RuGe_{12}^-$, the intensity slowly increases from 2.60 to 3.81 eV and the VDE is around 3.81 eV.

Theoretical results. The optimized global minima of all $RuGe_n^-$ and $RuGe_n$ ($n = 3-12$) clusters as obtained at the PW91/LANL2DZ level are plotted in Fig. 2. The geometries and relative stabilities of the low-lying isomers of each cluster are shown in Figs S1 and S2 in the Supporting Information (SI). The electronic states, point group symmetries and relative energies of all the considered clusters are summarized in Table S1 in the SI. The theoretical VDE and ADE values of the global minimum anions are listed in Table 1, in comparison with the experimental data. From Table 1, for all the ground state species, the computational VDE values are in good agreement with the experimental data, lending considerable credence to their structures.

As shown in Fig. 2, the global minimum structures of the neutral $RuGe_n$ clusters are nearly identical to those of $RuGe_n^-$ when $n = 6, 7, 9, 10, 11$, while the structures of $RuGe_n$ ($n = 3-5, 9, 12$) are different from their corresponding anions. However, both the anionic and neutral $RuGe_n^{-/0}$ clusters display the following structural trends: the global minimum $RuGe_n^{-/0}$ clusters with $n = 3-6$ have exohedral or Ru-capped geometries (except for $RuGe_5^-$); $RuGe_7^{-/0}$ and $RuGe_8^{-/0}$ possess half-encapsulated structures; and endohedral geometries are adopted for the larger clusters with $n = 9-12$. In particular, the global minimum structure of $RuGe_{10}^{-/0}$ is a 3-connected C_s-symmetric polyhedral cage encapsulating the Ru atom in the center. As well as, a C_2-symmetric ground state
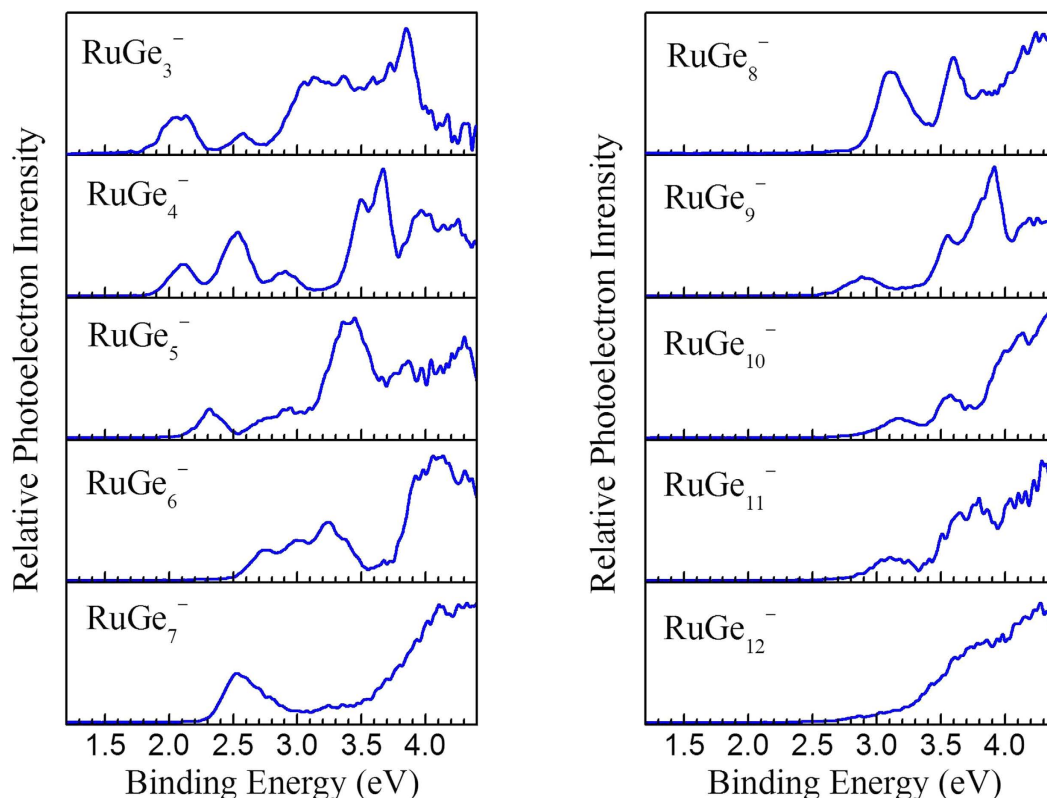
Figure 1. Experimental photoelectron spectra of RuGe$_n^-$ ($n = 3$–12) clusters recorded with 266 nm photons.

Cluster	VDE (eV)		ADE (eV)	
	Exp.	Theo.	Exp.	Theo.
RuGe$_3^-$	2.08 ± 0.08	1.98	1.80 ± 0.08	1.97
RuGe$_4^-$	2.12 ± 0.08	2.08	1.83 ± 0.08	2.37
RuGe$_5^-$	2.32 ± 0.08	2.64	2.08 ± 0.08	2.41
RuGe$_6^-$	2.75 ± 0.08	2.64	2.45 ± 0.08	2.77
RuGe$_7^-$	2.53 ± 0.08	2.58	2.27 ± 0.08	2.49
RuGe$_8^-$	3.10 ± 0.08	2.97	2.89 ± 0.08	2.68
RuGe$_9^-$	2.89 ± 0.08	3.17	2.53 ± 0.08	2.78
RuGe$_{10}^-$	3.17 ± 0.08	3.33	2.73 ± 0.08	3.11
RuGe$_{11}^-$	3.14 ± 0.08	3.19	2.79 ± 0.08	2.96
RuGe$_{12}^-$	3.81 ± 0.08	3.62	3.11 ± 0.08	3.41

Table 1. Vertical detachment energies (VDEs) and adiabatic detachment energies (ADEs) of the ground state RuGe$_n^-$ ($n = 3$–12) clusters estimated from their photoelectron spectra.

RuGe$_{12}^-$ also reveals a 3-connected polyhedral geometry, which mirrors the structure of trianionic RuGe$_{12}^{3-}$ reported by Espinoza-Quintero et al.[30] Unfortunately, except for this, there are no more available experimental data to compare with our calculations for RuGe$_n^{-/0}$.

Comparison between experiment and theory.   In order to prove the credibility of the obtained ground state structures, the photoelectron spectra of RuGe$_n^-$ ($n = 3$–12) are simulated and displayed in Fig. 3, along with the experiment spectra from Fig. 1 for comparison. In general, the simulated spectra are in overall satisfying agreement with the measured photoelectron spectra, which certify the validity of the present theoretical results.

For RuGe$_3^-$, the first two discrete peaks of the simulated spectrum are red-shifted compared to the experimental results, however, the others agree well. In the simulated spectrum of RuGe$_4^-$, the first five peaks are all in excellent agreement with the experimental measurements. For RuGe$_5^-$, the first peak is located at 2.64 eV, a little higher than the experimental value. The following peak is very broad, in accordance with experimental determination, and originates from several individual excitations. In the case of RuGe$_6^-$, the essential features of the experimental spectrum are well reproduced by DFT simulation. In the calculated photoelectron spectrum of
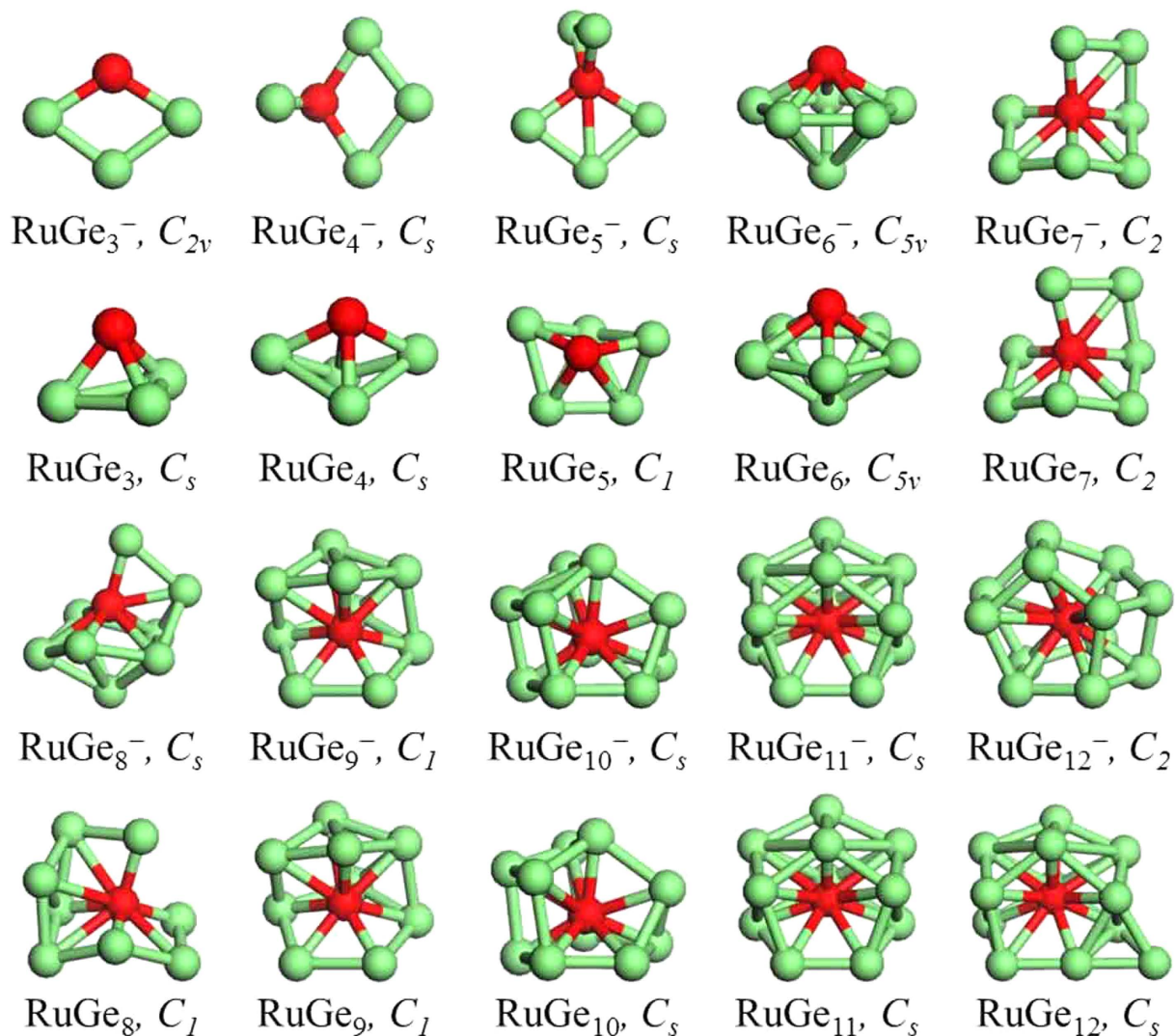
Figure 2. Global minimum structures of the RuGe$_n^-$ and RuGe$_n$ ($n = 3$–12), along with the point group symmetries. The green balls are Ge atoms and the red balls are Ru atoms.

RuGe$_7^-$, the first peak is in excellent agreement with the experimental measurement, while the others are somewhat different. In the theoretical spectrum of RuGe$_8^-$, three major peaks are obtained, which are in very good agreement with experiment. In the case of RuGe$_9^-$, the overall experimental spectrum is well reproduced by theoretical calculations; however, the first calculated peak at 3.17 eV is a little higher in energy and thus less distinct than the first experimental peak. In the theoretical spectrum of RuGe$_{10}^-$, three obvious peaks can be clearly seen, in accordance with the experimental data. For RuGe$_{11}^-$, the DFT calculation successfully reproduces the experimental trend and yields two obvious peaks. The simulated spectrum of RuGe$_{12}^-$ has an onset very similar to the experimental spectrum, and reproduces a small shoulder before the main peak.

For the sake of a clearer comparison to experiment, the experimental and computed VDEs and ADEs as functions of cluster size are shown in Fig. 4. In the experimental VDE curve, an overall increase can be seen from 1.98 to 3.62 eV across the different clusters, but features local minima at RuGe$_7^-$, RuGe$_9^-$, and RuGe$_{11}^-$. The theoretical VDEs reproduce this trend generally well, bar for the local minimum at RuGe$_9^-$. In general, the ADE value of an anionic cluster is equal to the adiabatic electron affinity (AEA) of the corresponding neutral species when their geometries are similar to each other. From experiment, it can be seen that ADE values keep rising from 1.80 to 2.45 eV for RuGe$_n^-$ ($n = 3$–6); thereafter this trend slows down, with an increase from 2.53 to 3.11 eV as the number of germanium atoms increases from 9 to 12. The ADE shows oscillatory from $n = 6$ to 9, with local minimum at RuGe$_7^-$ (2.27 ± 0.08 eV) and RuGe$_9^-$ (2.53 ± 0.08 eV). Again, the experimental trend is well reproduced by DFT simulations, but with an overestimation of about 0.46 eV at RuGe$_4^-$ and 0.30 eV at RuGe$_{10}^-$ respectively. From electronic points of view, RuGe$_7^-$ and RuGe$_9^-$ stand out as unique species with appreciable stability. For a given neutral cluster, lower AEA corresponds to higher stability. This indicates, in other words, that the neutral RuGe$_7$ and RuGe$_9$ clusters are stable.
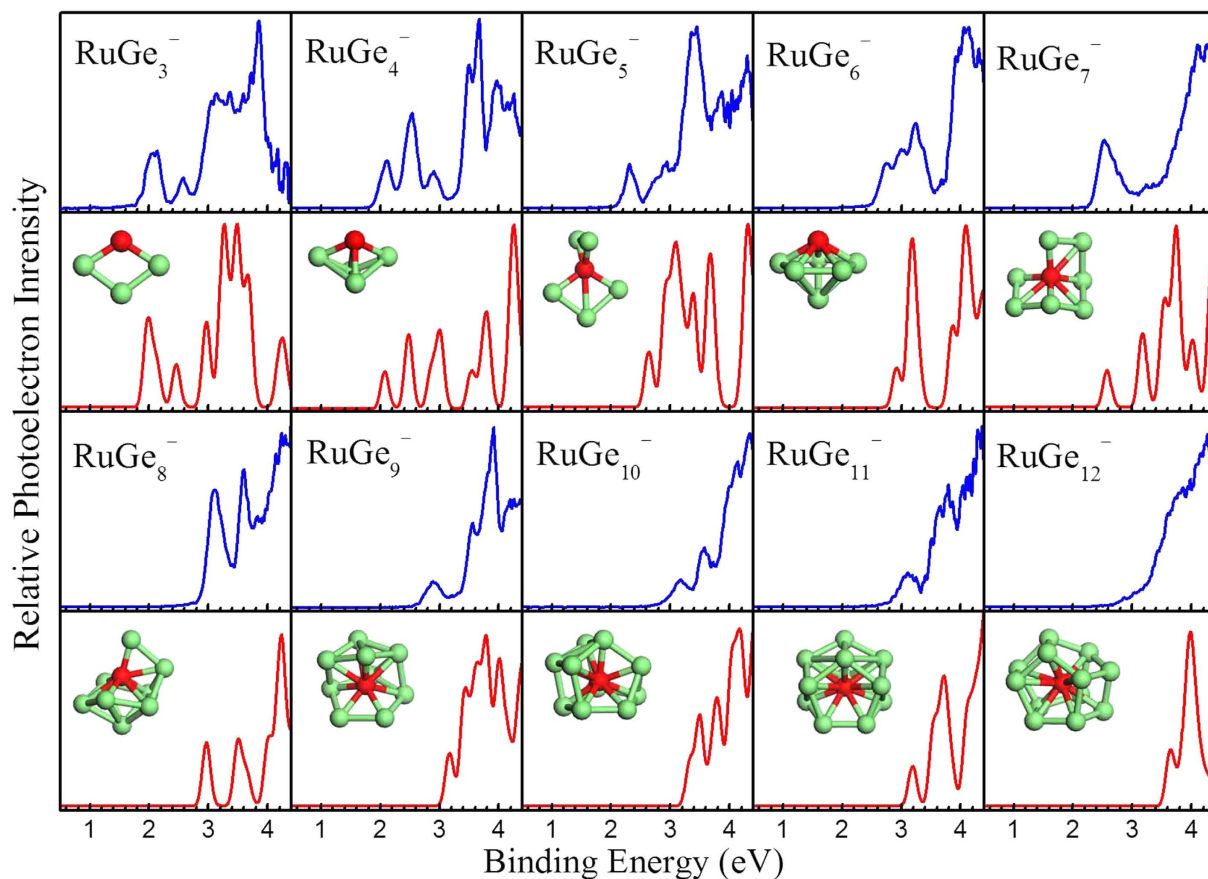
Figure 3. Simulated (red color) photoelectron spectra for RuGe$_n^-$ ($n = 3$–12) clusters, along with the experimental spectra (blue color) from Fig. 1 for comparison. A uniform Gaussian broadening of 0.15 eV is chosen for all the simulated spectra.
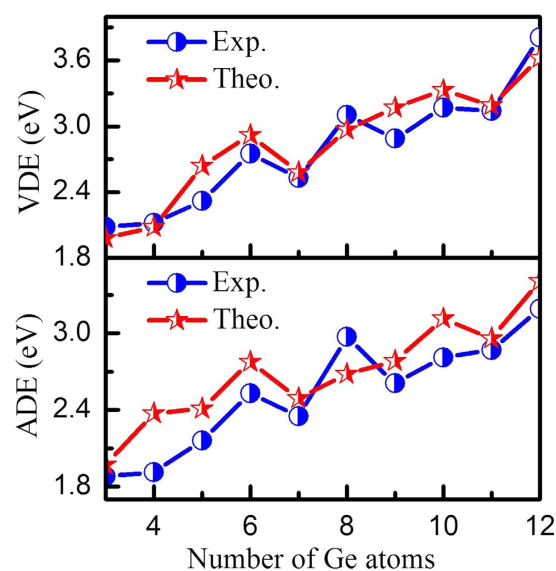

Figure 4. Vertical detachment energies (VDEs) and adiabatic detachment energies (ADEs) of RuGe$_n^{-/0}$ ($n = 3$–12) clusters: blue circles, experiment; red pentacles, theory.

Stabilities and electronic properties. The thermodynamic stabilities of the RuGe$_n^-$ and RuGe$_n$ ($n = 3$–12) clusters can be explored by investigating two thermodynamic parameters, namely the average binding
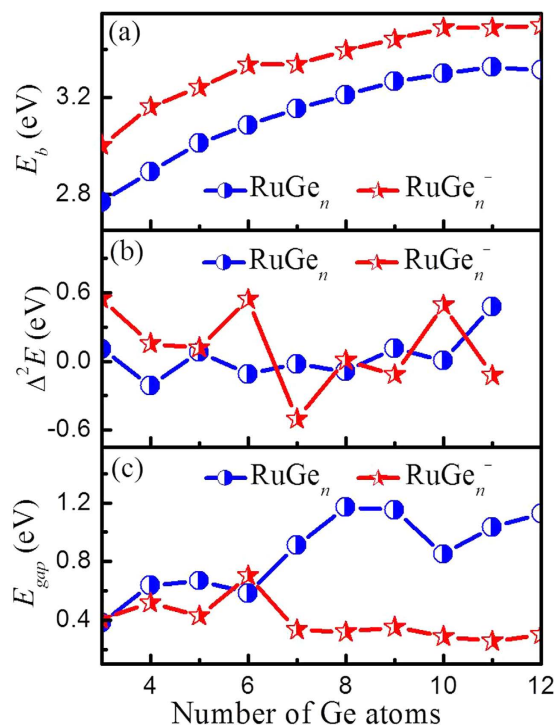
Figure 5. Size dependences of average binding energies (E_b), second order difference ($\Delta^2 E$) and HOMO–LUMO gaps (E_{gap}) for the global minimum $RuGe_n^{-/0}$ ($n = 3–12$) clusters.

energy and the second energy difference. The E_b and $\Delta^2 E$ of $RuGe_n^{-/0}$ clusters are defined as follows:

$$E_b(RuGe_n) = [E(Ru) + nE(Ge) - E(RuGe_n)]/(n+1) \tag{1}$$

$$E_b(RuGe_n^-) = [E(Ru) + E(Ge^-) + (n-1)E(Ge) - E(RuGe_n^-)]/(n+1) \tag{2}$$

$$\Delta^2 E(RuGe_n^{-/0}) = E(RuGe_{n-1}^{-/0}) + E(RuGe_{n+1}^{-/0}) - 2E(RuGe_n^{-/0}) \tag{3}$$

where E is the energy of the corresponding atom or cluster. The behaviors of calculated E_b and $\Delta^2 E$ values as function of cluster size for $RuGe_n^-$ and $RuGe_n$ clusters are shown in Fig. 5(a,b), respectively. From Fig. 5(a), the binding energy per atom increases monotonously, but saturates for $n > 9$, with increasing number of Ge atoms, suggesting that the formation of larger clusters is favorable over smaller clusters. For all cluster sizes, E_b values of anionic clusters are higher than those of the neutral species. This implies that as soon as a neutral cluster gains an extra electron, it becomes more stable. From Fig. 5(b), one can see stronger oscillations of the second energy difference in the anionic than the neutral clusters. An obvious odd−even oscillation for anions is found from $n = 5$ to 11. The peaks for $RuGe_6^-$, $RuGe_8^-$ and $RuGe_{10}^-$ indicate that these are more stable than their neighboring sized clusters.

The energy gap (E_{gap}) between the highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO) indicates the ability of electrons to jump from occupied orbitals to unoccupied orbitals. As shown in Table S1 and Fig. 5(c), the HOMO–LUMO gap values of $RuGe_n^{-/0}$ clusters range between 0.26 and 1.17 eV. Regarding cluster size, the E_{gap} values of neutral species with $n > 6$ are much larger than those of the small clusters. In contrast, the E_{gap} values of the larger sized anions are smaller than those of the small species. It stands out that neutral $RuGe_8$, $RuGe_9$ and $RuGe_{12}$ have the largest HOMO–LUMO gaps, implying that they are (relatively) more stable based on their electronic structure.

In order to examine the charge transfers, we next conduct the natural population analyses for the most stable isomers of the anionic and neutral $RuGe_n^{-/0}$ ($n = 3–12$) clusters. The results are summarized in Fig. 6. For the anionic clusters, it is found that, for $RuGe_3^-$, the negative charge is localized on both Ru atom and the Ge_n framework. For the larger anions, there is increasing electron transfer from the Ge_n framework to the Ru atom. For cluster size of $n = 7–12$, in particular, the negative charge on the Ru atom increases significantly, suggesting that there is substantial electron transfer from the Ge_n framework to the Ru atom. In the case of neutral clusters, the electron transfer from the Ge_n atoms to the Ru atom occurs for all cluster size and the charges on Ru are much larger for clusters with $n = 7–12$. The electron transfer from the Ge_n framework to the Ru atom is related to the formation of endohedral structures.
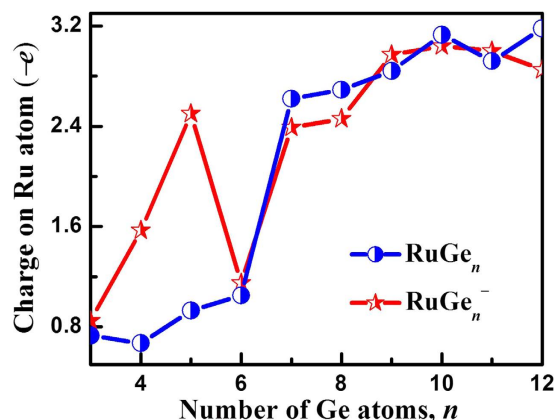
Figure 6. Natural charge populations of the Ru atom for the global minimum structures of $RuGe_n^{-/0}$ ($n = 3$–12) clusters.

Conclusions

In summary, we have reported a systematic study of the relative stability and electronic properties of ruthenium doped germanium clusters in the size range of $3 \leq n \leq 12$. Photoelectron spectra of anionic $RuGe_n^-$ clusters are measured at 266 nm. Unbiased structure searches reveal quite similar global minimum structures for both anionic and neutral clusters. Excellent agreement between theoretical calculations and experimental measurements is found. The global minimum anionic $RuGe_n^-$ clusters with $n = 3$–6 as well as their neutral counterparts have Ru-capped structures (except for $RuGe_5^-$); half-encapsulated structures are found for $RuGe_7^{-/0}$ and $RuGe_8^{-/0}$; and the larger clusters $RuGe_n^{-/0}$ ($n = 9$–12) feature endohedral geometries. From natural population analyses we see that, for the larger clusters with $n = 7$–12, the negative charge on the Ru atom increases significantly, suggesting that there is more electron transfer from the Ge_n framework to the Ru atom, which stabilizes the formation of endohedral cage.

Experimental and computational methods. Experiments are carried out on a home-built instrument consisting of a laser vaporization cluster source, a time-of-flight (TOF) mass spectrometer, and a magnetic-bottle photoelectron spectrometer, which have been described elsewhere[32]. In the laser vaporization source, the anionic Ru−Ge clusters are produced by laser ablation of a rotating translating disk target (13 mm diameter, Ru/Ge mole ratio 1:2) with a nanosecond Nd:YAG laser (Continuum Surelite II-10). To cool the generated cluster, helium gas with ~4 atm backing pressure is allowed to expand through a pulsed valve (General Valve Series 9) into the source. At the TOF mass spectrometer, the thus-formed cluster anions are mass-analyzed. The $RuGe_n^-$ ($n = 3$–12) clusters are individually selected with a mass gate, decelerated by a momentum decelerator, and crossed with the beam of a Nd:YAG laser (Continuum Surelite II-10, 266 nm) in the photodetachment setup. The magnetic-bottle photoelectron spectrometer is used to energy-analyze the electrons from the photodetachment. The resolution of the magnetic-bottle photoelectron spectrometer is about 40 meV at electron kinetic energy of 1 eV.

Theoretically, the structure searches for anionic $RuGe_n^-$ ($n = 3$–12) clusters and their neutral states are performed using CALYPSO method[33–37]. This method is based on globally minimizing potential energy surfaces, merging *ab initio* total energy calculations with CALYPSO cluster prediction through particle swarm optimization. It has been successful in correctly predicting structures for various systems[35,36,38]. The low-lying isomers of $RuGe_n^{-/0}$ ($n = 3$–12) found in the searches are further optimized using DFT with the PW91 functional[39]. The LANL2DZ basis set is used for both Ru and Ge atoms. Spin multiplicities (up to septet and octet) are considered for refined structure optimization. Vibrational frequency calculations are used to verify the nature of real local minima. To further evaluate the relative energies of the low-lying structures, single-point calculations were carried out by employing the larger def2-TZVP basis set. Excitation energies of the neutral cluster are calculated using the time-dependent DFT (TDDFT) method at the corresponding anionic structure[40]. All calculations are carried out with the Gaussian 09 package[41]. The PW91/LANL2DZ theory show superior results in terms of structures and binding energies for the species considered here and are therefore used for direct comparison with the experiment.

References

1. Liu, Y. R. *et al.* Three-dimensional assignment of the structures of atomic clusters: An example of Au₈M (M = Si, Ge, Sn) anion clusters. *Sci Rep-uk.* **5,** 17738 (2015).
2. Baletto, F. & Ferrando, R. Structural properties of nanoclusters: Energetic, thermodynamic, and kinetic effects. *Rev. Mod. Phys.* **77,** 371–423 (2005).
3. Ding, L. P. *et al.* Understanding the structural transformation, stability of medium-sized neutral and charged silicon clusters. *Sci Rep-uk.* **5,** 15951 (2015).
4. Yoo, S. & Zeng, X. C. Structures and stability of medium-sized silicon clusters III. Reexamination of motif transition in growth pattern from Si₁₅ to Si₂₀. *J. Chem. Phys.* **123,** 164303 (2005).
5. Yoo, S. & Zeng, X. C. Structures and relative stabilities of medium-sized silicon clusters. IV. Motif-based low-lying clusters Si₂₁-Si₃₀. *J. Chem. Phys.* **124,** 054304 (2006).
6. Yoo, S., Shao, N., Koehler, C., Frauenheim, Th. & Zeng, X. C. Structures and relative stability of medium-sized silicon clusters V. Low-lying endohedral fullerene-like clusters Si₃₁-Si₄₀ and Si₄₅. *J. Chem. Phys.* **124,** 164311 (2006).

7. Xu, H. G., Kong, X. Y., Deng, X. J., Zhang, Z. G. & Zheng, W. J. Smallest fullerene-like silicon cage stabilized by a V_2 unit. *J. Chem. Phys.* **140,** 024308 (2014).
8. Huang, X. *et al.* Structures and electronic properties of $V_3Si_n^-$ ($n = 3–14$) clusters: A combined ab initio and experimental study. *J. Phys. Chem. C* **119,** 10987–10994 (2015).
9. Wang, L. S., Wu, H., Desai, S. R., Fan, J. & Colson, S. D. A photoelectron spectroscopic study of small silicon oxide clusters: SiO_2, Si_2O_3, and Si_2O_4. *J. Phys. Chem.* **100,** 8697–8700 (1996).
10. Grubisic, A., Wang, H. P. & Ko, Y. Photoelectron spectroscopy of europium-silicon cluster anions, $EuSi_n^-$ ($3 \leq n \leq 17$). *J. Chem. Phys.* **129,** 054302 (2008).
11. Deng, X. J. *et al.* Photoelectron spectroscopy and density functional calculations of VGe_n^- ($n = 3–12$) clusters. *J. Phys. Chem. C* **119,** 11048–11055 (2015).
12. Deng, X. J., Kong, X. Y., Xu, X. L. & Xu, H. G. Structural and magnetic properties of $CoGe_n^-$ ($n = 2–11$) clusters: Photoelectron spectroscopy and density functional calculations. *ChemPhysChem* **15,** 3987–3993 (2014).
13. Palyanov, Y. N., Kupriyanov, I. N., Borzdov, Y. M. & Surovtsev, N. V. Germanium: A new catalyst for diamond synthesis and a new optically active impurity in diamond. *Sci Rep-uk.* **5,** 14789 (2015).
14. Pillarisetty, R. Academic and industry research progress in germanium nanodevices. *Nature* **479,** 324−328 (2011).
15. Wang, L. & Zhao, J. Competition between supercluster and stuffed cage structures in medium-sized Ge_n ($n = 30–39$) clusters. *J. Chem. Phys.* **128,** 024302 (2008).
16. Wang, J., Wang, G. & Zhao, J. Structure and electronic properties of Ge_n ($n = 2–25$) clusters from density-functional theory. *Phys. Rev. B.* **64,** 205411 (2001).
17. Wang, J. Q., Stegmaier, S. & Fässler, T. F. $[Co@Ge_{10}]^{3-}$: An intermetalloid cluster with archimedean pentagonal prismatic structure. *Angew. Chem. Int. Ed.* **48,** 2032–2036 (2009).
18. Zhou, B., Denning, M. S., Kays, D. L. & Goicoechea, J. M. Synthesis and isolation of $[Fe@Ge_{10}]^{3-}$: A pentagonal prismatic zintl ion cage encapsulating an interstitial iron atom. *J. Am. Chem. Soc.* **131,** 2802–2803 (2009).
19. Kumar, V. & Kawazoe, Y. Metal-encapsulated caged clusters of germanium with large gaps and different growth behavior than silicon. *Phys. Rev. Lett.* **88,** 235504 (2002).
20. Tai, T. B. & Nguyen, M. T. Enhanced stability by three-dimensional aromaticity of endohedrally doped clusters $X_{10}M^{0/-}$ with X = Ge, Sn, Pb and M = Cu, Ag, Au. *J. Phys. Chem. A* **115,** 9993–9999 (2011).
21. Wang, J. & Han, J. G. Geometries and electronic properties of the tungsten-doped germanium clusters: WGe_n ($n = 1−17$). *J. Phys. Chem. A* **110,** 12670–12677 (2006).
22. Bandyopadhyay, D. & Sen, P. Density functional investigation of structure and stability of Ge_n and Ge_nNi ($n = 1−20$) clusters: Validity of the electron counting rule. *J. Phys. Chem. A* **114,** 1835–1842 (2010).
23. Jing, Q., Tian, F. & Wang, Y. No quenching of magnetic moment for the Ge_nCo ($n = 1−13$) clusters: First-principles calculations. *J. Chem. Phys.* **128,** 124319 (2008).
24. Wang, J. & Han, J. G. A computational investigation of copper-doped germanium and germanium clusters by the density-functional theory. *J. Chem. Phys.* **123,** 244303 (2005).
25. Kumar, V., Singh, A. K. & Kawazoe, Y. Smallest magic caged clusters of Si, Ge, Sn, and Pb by encapsulation of transition metal atom. *Nano Lett.* **4,** 677−681 (2004).
26. Lu, J. & Nagase, S. Metal-doped germanium clusters MGe_n at the sizes of $n = 12$ and 10: Divergence of growth patterns from the MSi_n clusters. *Chem. Phys. Lett.* **372,** 394−398 (2003).
27. Claes, P. *et al.* Structural identification of caged vanadium doped silicon clusters. *Phys. Rev. Lett.* **107,** 173401 (2011).
28. Kawamura, H., Kumar, V. & Kawazoe, Y. Growth, magic behavior, and electronic and vibrational properties of Cr-doped Si clusters. *Phys. Rev. B* **70,** 245433 (2004).
29. Kawamura, H., Kumar, V. & Kawazoe, Y. Growth behavior of metal-doped silicon clusters Si_nM (M = Ti, Zr, Hf; $n = 8–16$). *Phys. Rev. B* **71,** 075423 (2005).
30. Espinoza-Quintero, G., Duckworth, J. C. A., Myers, W. K., McGrady, J. E. & Goicoechea, J. M. Synthesis and characterization of $[Ru@Ge_{12}]^{3-}$: An endohedral 3-connected cluster. *J. Am. Chem. Soc.* **136,** 1210–1213 (2014).
31. Goicoechea, J. M. & McGrady, J. E. On the structural landscape in endohedral silicon and germanium clusters, $M@Si_{12}$ and $M@Ge_{12}$. *Dalton Trans.* **44,** 6755–6766 (2015).
32. Xu, H. G. *et al.* Vanadium-doped small silicon clusters: Photoelectron spectroscopy and density-functional calculations. *Chem. Phys. Lett.* **487,** 204–208 (2010).
33. Wang, Y., Lv, J., Zhu, L. & Ma, Y. Crystal structure prediction via particle-swarm optimization. *Phys. Rev. B* **82,** 094116 (2010).
34. Wang, Y., Lv, J., Zhu, L. & Ma, Y. CALYPSO: A method for crystal structure prediction. *Comput. Phys. Commun.* 183, 2063–2070 (2012).
35. Lv, J., Wang, Y., Zhu, L. & Ma, Y. Particle-swarm structure prediction on clusters. *J. Chem. Phys.* **137,** 084104 (2012).
36. Lu, S., Wang, Y., Liu, H., Miao, M. & Ma, Y. Self-assembled ultrathin nanotubes on diamond (100) surface. *Nature Commun.* **5,** 3666 (2014).
37. Wang, Y. *et al.* An effective structure prediction method for layered materials based on 2D particle swarm optimization algorithm. *J. Chem. Phys.* **137,** 224108 (2012).
38. Zhu, L., Liu, H., Pickard, C. J., Zou, G. & Ma, Y. Reactions of xenon with iron and nickel are predicted in the Earth's inner core. *Nature Chem.* **6,** 644–648 (2014).
39. Perdew, J. P. & Wang, Y. Accurate and simple analytic representation of the electron-gas correlation energy. *Phys. Rev. B* **45,** 13244–13249 (1992).
40. Van Gisbergen, S. J. A., Snijders, J. G. & Baerends, E. J. Implementation of time-dependent density functional response equations. *Comput. Phys. Commun.* **118,** 119–138 (1999).
41. Frisch, M. J. *et al.* Gaussian 09, revision C.0; Gaussian, Inc.: Wallingford, CT (2009).

Acknowledgements

Author Contributions

X.K. and C.L. conceived the idea. Y.J., C.L. and C.Z. performed the calculations. S.L., H.X. and W.Z. performed the experiment. Y.J., A.H., C.L. and W.Z. wrote the manuscript and all authors contributed to revisions.

Additional Information

Supplementary information accompanies this paper at http://www.nature.com/srep